

Output Stages

Matthew Spencer

Harvey Mudd College

E151 – Analog Circuit Design

1

In this video series we're going to discuss output stages, which are amplifiers that can drive high currents or small loads. As the name suggests, they're often used as the outputs of op-amps. Op-amps have to drive all sorts of weird loads, so a good output stage is an important key to making an op-amp useful. Next video series we're putting this together with differential amplifiers to build our first op-amp.

Why Use Output Stages and Max Power Transfer

Matthew Spencer

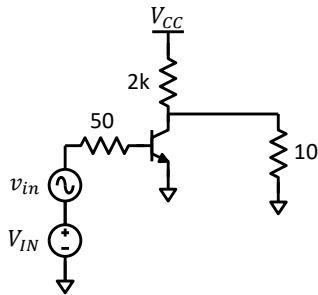
Harvey Mudd College

E151 – Analog Circuit Design

2

We haven't used output stages yet this semester, so why start now? This video will go over a quick justification for output stages, which are also sometimes called power amplifiers, and then cover an important theorem that lends us insight about how to transfer power into a load. Appropriately, the theorem is called the max power transfer theorem.

Small Loads Affect Bias Points and Gain



- Maximum $V_{OUT} = V_{CC} \cdot \frac{10\Omega}{2000k\Omega + 10\Omega}$

- What about AC coupling?

- Independent DC bias point,
- Small load sets $a_v = g_m R_L$.
- Coupling cap interferes w/ power Xfer

- Many small/tricky loads in the world: motors, speakers, plasmas, antennas, heaters, etc.

3

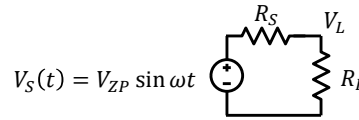
I've drawn a common emitter here that is DC coupled to a very small load resistance.

CLICK As you might have discovered in lab, very small load resistances mess with your bias point. In this case, we know that the highest voltage we could possibly get at the output node will occur when the transistor is in cutoff, and that output voltage still won't be very high. In fact, if that load is small enough, then the transistor might never be able to leave the saturation region.

CLICK OK, so what if we AC couple our loads. That's how we've been measuring output resistance, why wouldn't it work here? It would, in the sense that our collector's DC value would be whatever we had originally intended. However, the small load would still limit our gain, so it's not clear that we would be faithfully driving the v_{in} signal onto the load. Also, the coupling cap makes power transfer frequency dependent and difficult, which we will start to inspect on the next slide.

CLICK It's worth noting that small loads aren't an academic problem: there are lots of tiny, weird loads in the world like motors, speakers, plasmas, antennas heaters and others.

Max Power Transfer Theorem → Low r_{out}



Note: We're not analyzing complex Z_S and Z_L
 Note: This is a linear system, transistors aren't

Power in the Load:

$$P_L = \frac{V_L^2}{R_L} = \frac{1}{R_L} \left(\frac{R_L}{R_S + R_L} V_S \right)^2$$

$$P_L = \frac{R_L}{(R_S + R_L)^2} V_{ZP}^2 \sin^2 \omega t$$

$$\langle P_L \rangle = \frac{R_L}{(R_S + R_L)^2} \frac{V_{ZP}^2}{2}$$

Maximize in two cases – controlling R_L and controlling R_S :

$$\langle P_L \rangle = \frac{R_L}{(R_S + R_L)^2} \frac{V_{ZP}^2}{2} \rightarrow R_S = 0$$

$$\frac{d \langle P_L \rangle}{d R_L} = -3 \frac{V_{ZP}^2}{2} \left(\frac{1}{(R_S + R_L)^2} - \frac{2R_L}{(R_S + R_L)^3} \right)$$

$$= -3 \frac{V_{ZP}^2}{2} \left(\frac{R_S + R_L}{(R_S + R_L)^3} - \frac{2R_L}{(R_S + R_L)^3} \right)$$

$$\rightarrow R_L = R_S$$

Efficiency:

$$\eta = \frac{\langle P_L \rangle}{\langle P_S \rangle}$$

$$\eta_{MAX} = \frac{\left(\frac{V_S}{2}\right)^2 / R_L}{V_S^2 / R_L} = \frac{1}{4}$$

Because there are so many weird loads in the world, it's worth making a unified theory to figure out how to get power into them. That tool is called the maximum power transfer theorem. We derive it using the circuit pictured here.

CLICK Before we begin, I need to make two caveats. First, we're going to do this analysis using real resistors because that's the most familiar material in this class. However, there's a similar analysis that tells us a bit more information about reactive loads. Second, we're analyzing a linear system here, but we build amplifiers out of transistors, which are wildly non-linear. Even worse, we can't use our usual dodge of retreating to small signal analysis, because we want output stages to deliver high power, which means they need to have big voltage and current swings. However, you can still get a little bit of insight and/or ballpark estimates from small signal analysis at the extremes of an output stage's operation, and accordingly we can get a bit of insight or intuition from the max power transfer theorem.

CLICK We start by figuring out how much power is dissipated in the load. Because the load is resistive, that's given by the load voltage squared over R_L . The load voltage is expressed as a voltage divider multiplying V_S .

CLICK We can substitute the value for V_S and simplify some of the resistances out in front. Doing so reveals that the power is a function of time, specifically a sine squared.

CLICK We almost always care about the time averaged power instead of the instantaneous power. When we take a time average (which I've indicated with brackets), the sine squared turns into a factor of one half, which is a fact you should memorize.

CLICK Awesome, we have this expression for PL, let's see what we can do with it. We know that we want PL to be as big as possible, but there are two scenarios we have to consider. In the first, we control the source resistance. Just looking at the PL expression, we can see that RS only appears in the denominator, so we should minimize the source resistance to maximize the value of PL

CLICK The situation where we control RL is trickier to figure out. We have to take a derivative because RL is in both the numerator and denominator of our expression. I've done so here.

CLICK We can simplify that expression a bit in order to get everything over the same denominator. That means we can reduce our problem to finding RL such that these numerators, $RS+RL-2RL$, add to zero.

CLICK That's easy enough algebra, and it tells us that RL is equal to RS. Awesome, so now we know both how to send and receive power. Since we're building amplifiers, we're mostly in the case where we control the source, so the big insight we can get here is that we should design amplifiers that have very low output resistances. Now, the fact we're dealing with large signals complicates that message a bit since output resistance changes with bias point, but it's a good starting point for figuring out what an output stage might look like.

CLICK One final note, efficiency is a very fundamental measure of power transfer. It's the ratio of power delivered to the load divided by power pulled from the source. You might wonder what the maximum efficiency of matched power delivery looks like.

CLICK The answer is pretty bad. Because RS is equal to RL when we have maximum transfer, we know the voltage at the load is half the source voltage. Squaring that gives us an efficiency of $1/4$.

Summary

- We need special circuits to drive small loads, strange impedances, or lots of current. These are called power amplifiers or output stages.
- Maximize power transfer by minimizing R_S or matching R_L .
- Efficiency is $\eta = P_L/P_S$.
- High power circuits require large signal analysis.

Class A Amplifiers

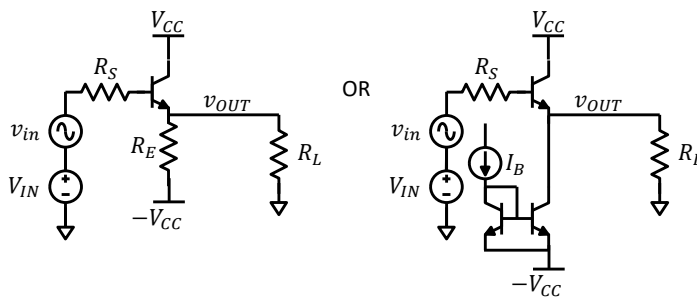
Matthew Spencer

Harvey Mudd College

E151 – Analog Circuit Design

In this video we're going to talk about the simplest type of output stage, which is called a class A amplifier.

Emitter Followers are Power Amplifiers



EF Increases Power Delivered:

$$v_{OUT} = v_{IN} - V_{BEON}$$

$$i_E = (\beta + 1)i_B$$

$$r_{out} \approx 1/g_m \text{ varies over cycle}$$

Resistive load calculations tricky:

$$v_{OUT} = v_{IN} - \phi_{TH} \ln \frac{i_C}{I_S} \quad \leftarrow \text{Transcendental}$$

Active Load Sets Max -ve I_L :

$$v_{OUT} = v_{IN} - \phi_{TH} \ln \frac{I_B + \frac{v_{OUT}}{R_L}}{I_S}$$

Class A amplifiers is a very fancy name that almost always just refers to an emitter follower. It has a few extra implications, but not that many. We'll talk about them on a later slide.

I've drawn two emitter followers here, one with a resistive bias and one with a current mirror bias. I've also chosen to power this emitter follower with split rails, both a positive VCC and a negative VCC. The negative VCC rail is often named VEE, but I've just put its voltage value there to simplify our notation. The split rails let us pick the DC bias VIN so that no power is dissipated in RL when the signal indicated by little vin is zero. I'm calling that signal little vin instead of small signal vin because it isn't really a small signal here. We're using small signal notation to indicate that little vin isn't constant, but it will swing basically rail to rail if we're trying to deliver power to RL.

CLICK Emitter followers have a small rout value, so we might have an instinct that they'll be good power amplifiers based on the max power transfer theorem. We need to do some large signal analysis to confirm that suspicion. First, we note that vOUT is just our total vIN signal minus VBEON. That means the amplitude of the voltage waveform on vOUT is the same as the one on vIN. However, the current driven into the output is going to be about beta times higher than the current driven into the base. That means we get the same voltage at a much higher current, which increases the power at our output. Finally, it's worth noting that the reason our initial suspicion wouldn't work is because rout changes as

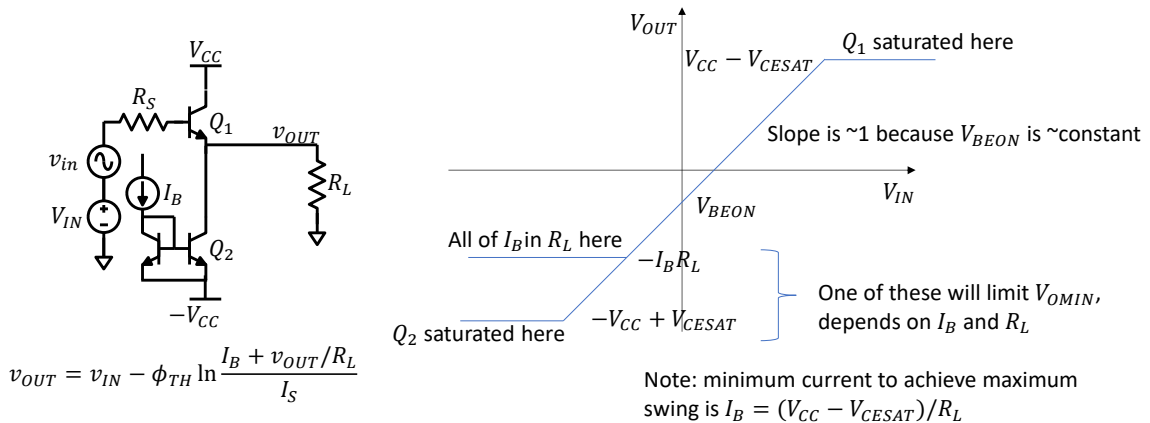
v_{IN} moves around because r_{out} is proportional to $1/g_m$ and g_m depends on i_E .

CLICK The resistively biased emitter follower is relatively unpopular because it's tough to analyze and, as a result, tough to design. You'll see something that looks like that circuit in a lot of places, notably in voltage regulators where v_{OUT} is controlled by feedback. (Also, voltage regulators take R_E out of the circuit and use R_L as the DC path to ground, which I find clever.) However, looking at a detailed calculation of the v_{OUT} voltage where V_{BEON} has been replaced by $\phi_{TH} \cdot \log(i_C/I_S)$, we realize this is going to be a brutal equation to solve. Notably, figuring out i_C is a transcendental resistor-diode problem.

CLICK The active load sidesteps that problem by pinning the total current that flows through the transistor in a convenient way. Finding v_{OUT} is still a transcendental equation, but we can make some observations about what happens at high and low extremes of v_{OUT} , notably that when v_{OUT}/R_L is equal to $-I_B$ we're going to see v_{BE} start to get really big.

... show R_L calculation with negative v_{IN} on resistive load?

Voltage Transfer Characteristic for DC V_{IN}



One easy way to start understanding the large signal behavior of this amplifier is drawing the voltage transfer characteristic or VTC, which you will recall is a plot of large signal V_{IN} vs. large signal V_{OUT} . I've put the axes over on the right to start doing that. I've also included the equation describing v_{OUT} from the last slide for reference.

CLICK We can start by noticing that the VTC will basically be a straight line for much of the curve. In this region I_B is bigger than v_{OUT}/R_L , so the logarithmic argument in the v_{OUT} equation is basically constant. This line has a value of zero when V_{IN} is equal to V_{BEON} . That's because of the V_{BEON} drop of Q_1 . Picking large signal $V_{IN}=V_{BEON}$ is a pretty common choice because then our output doesn't dissipate any power when little v_{in} is zero.

CLICK In the case that R_L is big, the v_{OUT}/R_L term never gets close to I_B , so the transfer characteristic has a slope very close to 1. We'll see that behavior change when the amplifier stops working because either Q_1 or Q_2 saturates, which happens when the output gets within V_{CESAT} of the rails. Hitting the Q_1 failure mode is unusual because it requires V_{IN} to swing above V_{CC} . The input is V_{BEON} above the output, so when the input is at V_{CC} , the output is at $V_{CC}-V_{BEON}$, which is still in the linear region of the amplifier's operation.

CLICK However, if R_L is not very large then there's another possible failure mode. If the current in Q1 goes to zero then it's cut off, so Q1 will be unable to convey additional changes in V_{IN} to the output. That happens at a voltage of $-I_B \cdot R_L$ because all of the current flowing into Q2 has been steered into R_L at that point. (You can imagine the purpose of Q1 in this circuit as stealing some of the Q2 current from R_L to increase the output voltage.) This causes the output voltage to be pinned at its lowest value, $-I_B \cdot R_L$. This appears in the v_{OUT} equation because the log function represents v_{BE} and the log function becomes very large as $-v_{OUT}/R_L$ cancels out I_B .

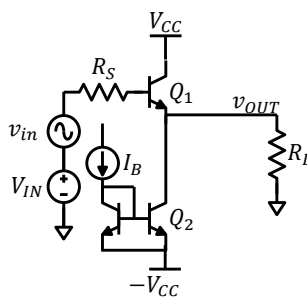
CLICK That means we have two possible conditions that set the minimum output voltage of the amplifier. Which one dominates depends on the values of I_B and R_L , which you can anticipate as a designer. One desirable condition is to have these two conditions occur at the same voltage so that you're neither constrained by your bias current, nor paying extra power for current you don't need. That occurs when $I_B \cdot R_L = 2(V_{CC} - V_{CESAT})$.

... $I_B \cdot R_L$ matches with V_{BEON} getting big. (Sat is different effect)

... Note that you need your input to swing outside the rails to reach Q1 saturation

Class A Efficiency Analysis

- Achieve max efficiency at maximum swing with minimum I_B . If you have too much I_B or too little swing then some I_B doesn't go into R_L .



$$\text{Let } I_B = 2(V_{CC} - V_{CESAT})/R_L$$

$$\text{Let } V_{IN} = 0.7V \text{ and } v_{in} = (V_{CC} - V_{CESAT}) \sin \omega t$$

Conditions for maximum η

$$P_S = P_{V_{CC}} + P_{-V_{CC}} = V_{CC}I_B(1 + \sin \omega t) + (-V_{CC})(-I_B)$$

$$\langle P_S \rangle = 2V_{CC}I_B$$

$$P_L = v_{OUT}i_L = (V_{CC} - V_{CESAT}) \sin \omega t \cdot I_B \sin \omega t$$

$$\langle P_L \rangle = \frac{1}{2}(V_{CC} - V_{CESAT})I_B$$

$$\eta_{MAX} = \frac{\langle P_L \rangle}{\langle P_S \rangle} = \frac{\frac{1}{2}(V_{CC} - V_{CESAT})}{2V_{CC}} = \frac{1}{4} \left(1 - \frac{V_{CESAT}}{2V_{CC}} \right) \approx 25\%$$

9

It's common to use efficiency as the main figure of merit for power amplifiers, so I want to show you one or two efficiency analyses in these videos to familiarize you with the how efficiency is calculated. However, we're not going to do a lot of efficiency analysis on our own in the lab, so enjoy the show and try to remember rough outlines of the process. You can just write down the results at the end of the page.

CLICK When you analyze efficiency, you usually do it in the most favorable condition for the amplifier. For this power amplifier, the most favorable condition is when the voltage swing is at its maximum value. That's because the I_B current is always flowing, whether you are driving a large output voltage or a small one. So we want to drive a big output voltage to use up all of the I_B current. For the same reason, we want to make sure I_B is as small as possible to achieve our maximum swing.

CLICK That gives rise to these test conditions. We're setting I_B to achieve maximum swing, which is a value we calculated on the last slide, we're setting large signal V_{IN} so that our output is zero centered, and we're picking little v_{in} to be a sinusoid with the maximum possible swing of $V_{CC} - V_{CESAT}$.

CLICK The power pulled from the supply is made up of power pulled from V_{CC} and power pulled into $-V_{CC}$. The power pulled into $-V_{CC}$ is easy: Q_2 is always sinking I_B of current.

The power from the positive supply is given by V_{CC} times I_C of Q1, which varies sinusoidally with the input. When v_{in} is zero, that current has to be I_B . When v_{in} goes to its max value, we pick up another $(V_{CC}-V_{CESAT})/R_L$ of current, which I've chosen to write as I_B here.

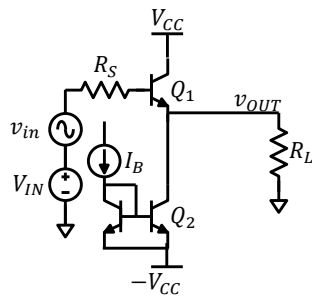
CLICK Time averaging this, we see that PS is $2 \cdot V_{CC} \cdot I_B$. This makes sense: on average I_B runs from $+V_{CC}$ to $-V_{CC}$.

CLICK The power delivered to the load is a bit more complicated because some power gets dissipated in Q1. The easiest way to figure it out here is multiplying v_{OUT} by i_{OUT} . We've set up our input voltages so that v_{OUT} is $V_{CC}-V_{CESAT}$ times sine omega t. Our current is just going to be proportional to that because our load is a resistor, but we've chosen to write the current coefficient as I_B instead of $(V_{CC}-V_{CESAT})/R_L$.

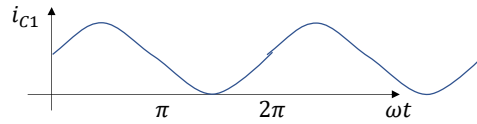
CLICK Time averaging the load power turns the sine squared into a factor of $1/2$.

CLICK Finally, we take the ratio of load power to source power to find our efficiency. The I_B values cancel out, which is why we kept writing stuff in terms of I_B above. We don't get particularly good news: the efficiency peaks at about 25%, which is quite bad! This is also our maximum efficiency, and it gets worse as we back off from maximum swing. The main cause of this is the fact we're always dissipated I_B from our supplies to keep Q1 in forward active.

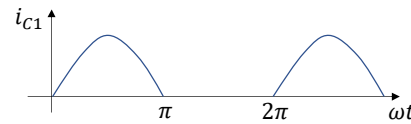
Conduction Angle Determines Class A/B/C



Class A:
 $\theta_c = 2\pi$

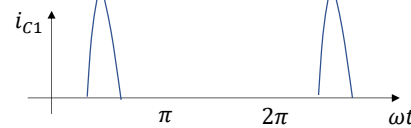


Class B:
 $\theta_c \approx \pi$



(same input as class A with smaller V_{IN})

Class C:
 $\theta_c < \pi$



(sketch presumes a very large drive amplitude and very -ve V_{IN})

So we might wonder if we can get away with not running current in Q1 all the time. That line of thinking leads us to some other very common amplifier topologies called class B and class C amplifiers. The circuit on the left can represent any of class A, B or C amplifiers, the only difference between them is the value of V_{IN} . More specifically, the way you determine if an amplifier is class A, B or C is a parameter called the conduction angle, which I've given the symbol capital ThetaC. The conduction angle describes the portion of the input wave's period during which Q1 is conducting current. In a class A amplifier, the conduction angle is 2π , which means it conducts current during the whole cycle of the input wave. You can see in the top-most plot of Q1's collector current vs. ωt , which is the usual way to look at conduction angle. Using ωt on the x axis let's us talk about fractions of a period without paying attention to frequency, which is where the concept of a conduction angle comes from. i_{C1} is always greater than zero, so its conduction angle is 2π . Class B amplifiers have conduction angles that are close to π , and class C amplifiers have conduction angles that are much less than π . Class C amplifiers are also usually driven very hard, which is why I've made the current spikes in the class C graph very tall.

You might look at the class B and class C amplifier outputs and think they look nothing like our input sine wave. How can those possibly be useful amplifiers if they throw away so much information about our wave? There are circuit tricks that let us use both class B and

class C amplifiers to replicate the input pretty faithfully, and these tricks include push-pull configurations and resonant output filters. More on those later.

... i_{C1} can't be negative, current only flows down through Q1, so you can think of Q1 canceling effect of I_B on RL

... class B and C have high distortion right now. Can use output filters. We're also going to see a class B variant that does much better.

Summary

- Emitter followers are power amplifiers with $V_{OMAX} = V_{CC} - V_{CESAT}$ and $V_{OMIN} = \max(-I_B R_L, -V_{CC} + V_{CESAT})$
- Operated as class A amplifiers (2π conduction angle) at maximum swing, class A amplifiers achieve 25% efficiency.
- Backoff degrades efficiency because I_B is running all the time.
- Class B and C amps are the same topology with less conduction angle.

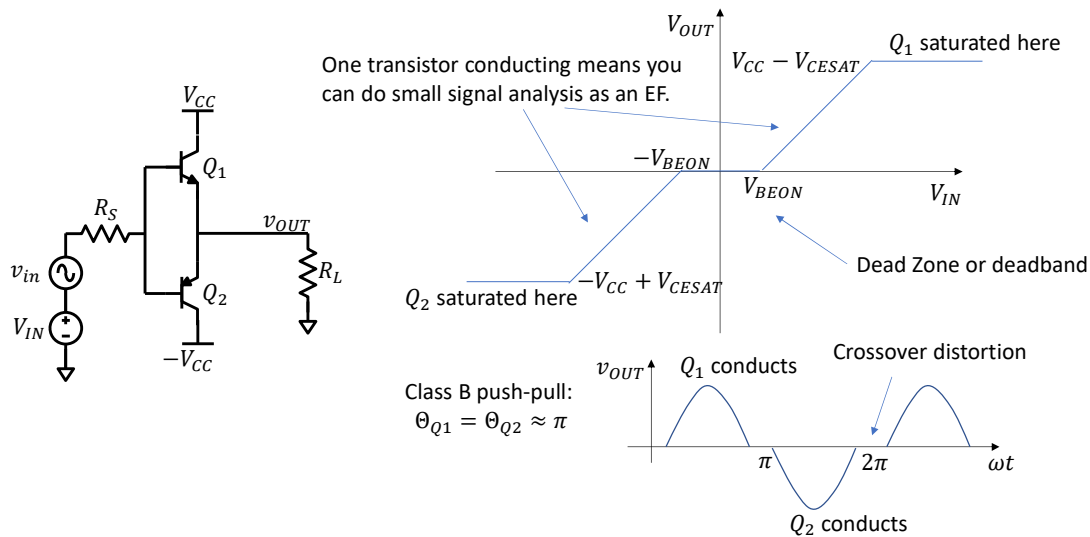
Class B and AB Push-Pulls

Matthew Spencer
Harvey Mudd College
E151 – Analog Circuit Design

12

In this video we're going to look at different types of power amplifiers called push-pulls. Push-pulls are very good amplifiers, and they are the most popular linear power amplifiers in use today.

Class B Push-Pulls Have Crossover Distortion



Our class B amplifier had a conduction angle of pi, which meant that we had really bad distortion at the output. One fix for that issue is to make a class B push-pull stage, which is pictured on the left. This kind of looks like two emitter followers, one NPN and one PNP, stacked on top of one another. Push pulls are set up to have Q1 conduct for pi radians while sourcing current to RL and then to have Q2 conduct for pi radians while sinking current from RL.

The axes for a transfer characteristic are set up pictured in the upper right.

CLICK We can start filling in the VTC by noting that when VIN is close to zero, neither Q1 or Q2 is turned on so VOUT doesn't change when VIN does. That behavior is called a dead zone or a deadband.

CLICK Once Q1 or Q2 turns on, that device acts like an emitter follower until it saturates. Just like class A amplifiers, reaching these saturation points it unusual because it requires VIN to swing outside of the rails.

CLICK One side note, if you need to do small signal analysis for some reason – maybe estimating peak output impedance or something – you can just treat the circuit as an emitter follower in these zones because one of the transistors is in cutoff. RL is going to act

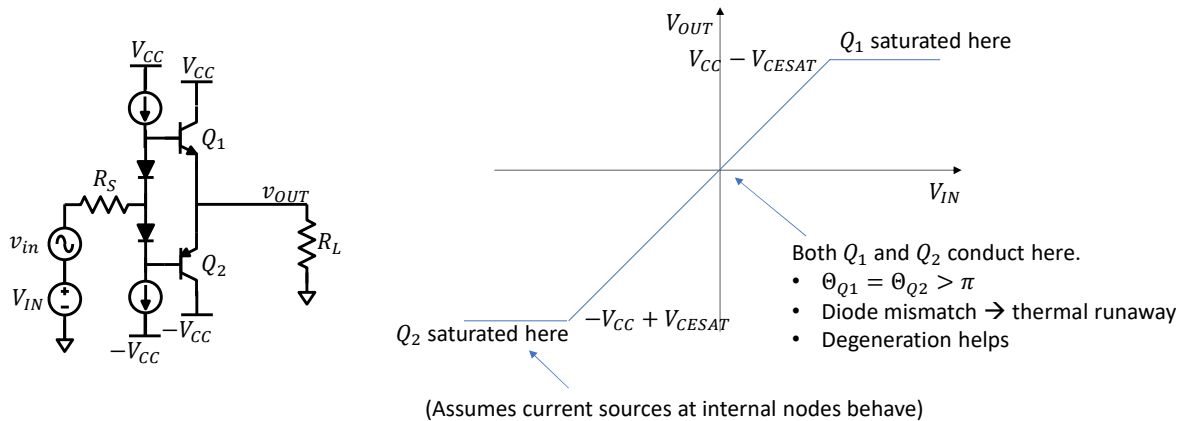
like the emitter resistor.

CLICK This is a weird voltage transfer characteristic, and it leads to a weird output wave. Specifically, our amplifier stops responding whenever the input crosses zero, which is called crossover distortion. That's bad, but you might not care if your input signal is sufficiently large. It's worth noting that both Q1 and Q2 have conduction angles that are close to π , which is why this called a class B push-pull. Q1 conducts during the positive half-cycle and Q2 conducts during the negative half-cycle.

... Crossover distortion

... Nice because biasing is easy

Class AB Fixes Crossover Distortion w/ Diodes



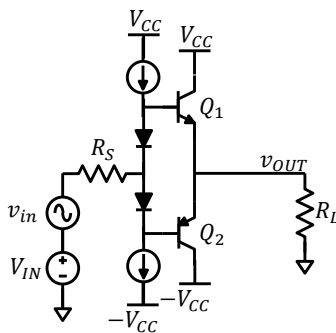
It's common to fix crossover distortion using a topology called a class AB push-pull. I've drawn one of those on the left. The big idea in a class AB push pull is that the on-voltage of the diodes cancels out the V_{BEON} voltages of Q_1 and Q_2 . So, for example, the base node of Q_1 has a voltage of V_{ON} on it when V_{IN} is zero, which lets Q_1 respond instantly if V_{IN} increases above zero. That gives rise to the voltage transfer characteristic I've drawn on the right.

CLICK There are a few details to note about this plot. First, pedantically, this Q_2 saturation condition depends on the base node being able to swing outside the rail, so the current sources that keep the diodes turned on need to be very well behaved.

CLICK Second, there's a narrow window in the middle of this curve where both Q_1 and Q_2 are conducting. By definition, that means both Q_1 and Q_2 have conduction angles greater than π , which is why this is called class AB, it's part way between the 2π conduction angle of class A and the π conduction angle of class B. It also means that we're susceptible to a cool circuit failure mode called thermal runaway. Thermal runaway comes from the fact that V_{BEON} is dependent on temperature, so if a slight mismatch in the diodes makes Q_1 take more of the current during this crossover, then Q_1 will get hotter than Q_2 , that means the v_{BE} of Q_1 will shrink, so that it takes even more current on the next cycle. Eventually Q_1 handles all of the output current and lights itself on fire. Fortunately, you can avoid this

by putting small resistors in series with the Q1 and Q2 emitters and with the biasing diodes.

Class AB Push-Pull Efficiency Analysis



Let $V_{IN} = 0V$ and $v_{in} = (V_{CC} - V_{CESAT}) \sin \omega t$

Note: $v_{OUT} = v_{IN}$

$$\langle I_{VCC} \rangle = \frac{1}{\pi} \int_0^{\pi} \frac{v_{OUT}}{R_L} d\omega t = \frac{V_{CC} - V_{CESAT}}{\pi R_L}$$

$$\langle P_S \rangle = P_{VCC} + P_{-VCC} = 2V_{CC} \langle I_{VCC} \rangle = \frac{2V_{CC}(V_{CC} - V_{CESAT})}{\pi R_L}$$

$$\langle P_L \rangle = \frac{\langle V_L \rangle^2}{R_L} = \frac{1}{2} \frac{(V_{CC} - V_{CESAT})^2}{R_L}$$

$$\eta_{MAX} = \frac{\langle P_L \rangle}{\langle P_S \rangle} = \frac{\pi}{4} \frac{V_{CC} - V_{CESAT}}{V_{CC}} \approx 78\%$$

Class B and class AB amplifiers have better efficiency than class A amplifiers, and we're going to go through a derivation to show that. Like the other efficiency analysis, we're mostly interested in the results and a chance to see the process. This analysis is for class AB amplifiers, but it's pretty accurate for class B amplifiers too.

CLICK We set large signal VIN to zero so that we don't dissipate DC power and we set little vin to $(V_{CC} - V_{CESAT}) \sin(\omega t)$ to ensure that we're driving this amplifier at maximum swing. We'll see soon that our efficiency is still best when our swing is as large as possible.

CLICK I want to pause for a moment to note that if everything is working right in this amplifier, then our output voltage is equal to our input voltage. The base of Q1 is at $v_{in} + V_{ON}$ and the output is back down to $v_{in} + V_{ON} - V_{BEON}$, and V_{ON} cancels out V_{BEON} to leave $v_{out} = v_{in}$.

CLICK We're going to care how much current is drawn from the supply, so I'm going to start by calculating it separately here. We're want to average our current, which is given by v_{OUT}/R_L , and we're going to do so using the conduction angle instead of absolute time. Q1 only conducts from 0 to pi (more or less), so we integrate v_{OUT} from 0 to pi with respect to ωt . Because we're finding an average, we need to divide by one over the period out in front of the integral. Integrating a half-period of sine gives us 1, so our output

is $(V_{CC} - V_{CESAT}) / (\pi * R_L)$.

CLICK The power pulled from the supply is the power pulled from V_{CC} plus the power sunk into $-V_{CC}$. The power pulled from V_{CC} is V_{CC} times the current we just calculated, and the current sunk by $-V_{CC}$ is an exactly symmetrical calculation using $-V_{CC}$ and the current through Q2 in the negative half-cycle. So our result is $2 * V_{CC} * (\text{the time average of } I_{CC})$ or this value shown on the right.

CLICK The power in the load is easier because it's a resistor. v_{OUT} squared has a sine squared in it, and time averaging that gives us a factor of $1/2$.

CLICK Finally, we find efficiency by taking the ratio of average load power to average source power. This comes out to $\pi/4$ times $(V_{CC} - V_{CESAT}) / V_{CC}$. If V_{CESAT} is small, this result is about 78%, which is more than three times better than our class A amplifier. We also back off more gracefully because our I_{VCC} scales with the input amplitude. However, that current is still pulled from a full V_{CC} , so we are most efficient when we are pulling as much current into the load as possible.

Summary

- Push-pulls use two devices, one to source current for about half the cycle and one to sink current for the other half.
- Class B push pulls have crossover distortion.
- Class AB push pulls used diodes to cancel crossover distortion, but have a risk of thermal runaway.
- Class B and class AB achieve peak efficiency of ~78%

Other Power Amplifiers

Matthew Spencer

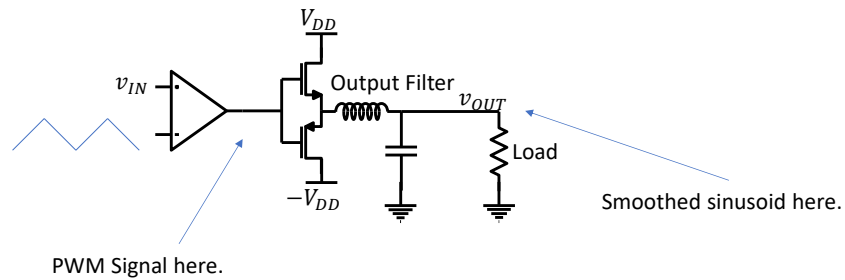
Harvey Mudd College

E151 – Analog Circuit Design

17

In this video we're going to do a quick survey of a few other power amplifiers. We're not analyzing these, I just want you to see some of the other technologies that are out there.

Class D Amplifiers Rely on PWM Signals

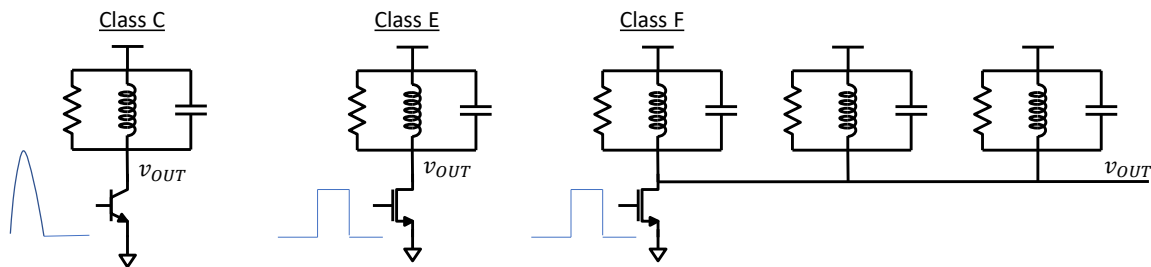


18

Class D amplifiers work quite differently from class A, B or C amplifiers. They work by driving a PWM signal onto a pair of switches, then filtering the output of those switches with a lossless LC output filter. That means the average power delivered to the output is going to depend on the PWM value, but the filter is going to force the PWM signal into a smooth shape that matches v_{IN} , usually a sinusoid. These are popular in audio applications because you can run the PWM at very high frequencies compared to the relatively low audio frequencies contained in the v_{IN} signal. However, hardcore audiophiles are skeptical of class D amplifiers because they claim they can hear faint artifacts of digital modulation.

By the way, the op-amp on the left of this figure shows a neat way to make a duty cycled output. By driving a triangle wave into one input of the op-amp and comparing it to our input signal, we get a square wave at the output with a width that is proportional to v_{IN} . That's because the time at which the triangle crosses past v_{IN} varies linearly with the v_{IN} value.

Class B, C, E, F Amps Rely on Output Filters



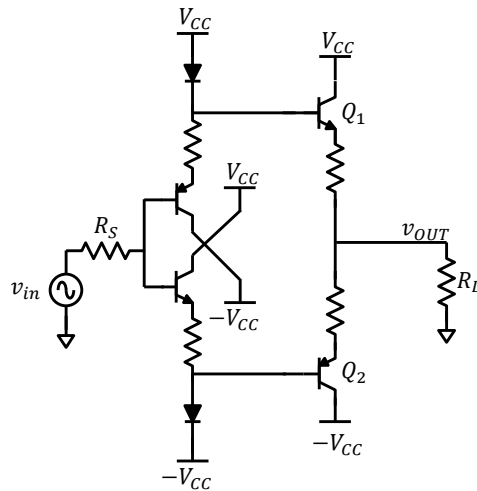
19

You might have noticed the output filter on the class D amplifier on the last slide and though back to our class B and C power amplifiers. Class B amplifiers can be made into push-pulls, but class C amplifiers had a very narrow conduction angle that guaranteed a weird output waveform. Using a resonant output filter can fix that, and that technique is the basis of class C, class E and class F amplifiers. The output filter turns a weird input wave into a sine wave at the vin frequency, and because it's lossless it boosts efficiency when it does so.

The difference between class C and class E amplifiers is that class C amplifiers try to use the conducting device in the linear region for at least a little while, while class E amplifiers are just switched digitally. Class F amplifiers add more output filters to this setup to try to suppress higher harmonics that sneak through the first filter. Each additional filter increases efficiency a bit by squeezing energy back into the vin frequency, but each adds a bit of loss due to parasitics. So there's a careful tradeoff to designing this big filter network.

These filters are popular at RF frequencies where passive components are small, so it's cheaper to build inductors and capacitors. It's also worth noting that these amplifiers are very efficient, and using a technique called zero voltage switching, where the digital changes are carefully timed with the output waveform, lets class E and class F amplifiers get into the 95% to 99% efficiency range.

Diamond Buffers are Ferocious Dinosaurs



20

Finally, I have to mention one of my favorite curiosities from prehistoric times, the diamond buffer. This is a class AB output stage where the diode drops that cancel the dead zone are provided by a pair of emitter followers. This circuit has very high input impedance and output drive current capacity, and I've never met a problem where it failed to drive a load. It also uses very high voltage rails, which means the biasing diodes on the emitter followers are often made with LEDs. These double as convenient debugging indicators: if one gets too bright then a component is going to start smoking soon. This is a fun one to have in your back pocket.

Summary

- There are a lot of fun power amplifiers in the world!